SHACK-HARTMANN TOMOGRAPHY FOR MULTIMODE OPTICAL BEAM PROPAGATION B. Stoklasa^{1,2}, L. Motka^{1,2}, J. Rehacek¹, Z. Hradil¹, L. L. Sanchez-Soto³ ¹Department of Optics, Palacky University, Olomouc, Czech Republic ²Meopta-Optika, Prerov, Czech Republic ³Departamento de Optica, Universidad Complutense, Madrid, Spain

Motivation

We experimentally show that wavefront detection combined with tomography processing can be used for the complete characterization of the second-order coherence and hence 3D imaging of partially coherent vortex beams. A standard Shack-Hartmann single-shot measurement is used in the method.



Experimental setups

To be able to propagate intensity, mutual coherence function *G* must be known at the input:

 $I(x) = \int_{-\infty}^{\infty} h(x, x') h^*(x, x'') G(x', x'') dx' dx''$ Let ρ be a coherence matrix, a discrete version of function G(x', x''):

 $G(x',x'') = \langle x'|\rho|x''\rangle = \operatorname{Tr}(\rho|x'\rangle\langle x''|)$

If basis for the problem is labeled by $|k\rangle$ (k = 1, ..., d, with d being the dimension), the complex amplitudes are $\langle x|k\rangle = \psi_k(x)$. Therefore, the coherence matrix ρ is given by $d \times d$ non-negative matrix.

Principle of SH tomography



If a SH sensor is illuminated with a coherent signal U(x), and the *i*th microlens is Δx_i apart from the SH axis, this microlens feels the field $U(x - \Delta x_i) = \langle x | \exp(-i\Delta x_i P) | U \rangle$. This field is truncated and filtered by the aperture (or pupil) function $A(x) = \langle x | A \rangle$ and Fourier transformed by the microlens:

 $U'(\Delta p_j) = \langle A | \exp(-i\Delta p_i X) \exp(-i\Delta x_i P) | U \rangle$. The intensity measured at the *j*th pixel behind the *i*th lens is then governed by a Born-like rule



Nd:YAG beam propagation

Shack-Hartmann tomography was composed of intensity data from 11x11 pixels against 7x7 microlenses, altogether 5929 measurements





 $I(\Delta x_i, \Delta p_j) = \operatorname{Tr}(\rho |\pi_{ij}\rangle \langle \pi_{ij}|),$ with $|\pi_{ij}\rangle = \exp(i\Delta x_i P) \exp(i\Delta p_i X) |A\rangle.$

Mode decomposition of SH tomography



We define the dynamical range (or field of view) of the SH tomography as the set of normal modes of measuring matrix with singular values exceeding a given threshold.

$$I_{ij} = \sum_{k} p_k^{ij} r_k$$





Indexes *i*, *j* label HG modes in the following order: $HG_{0,0}, HG_{1,0}, HG_{2,0}, HG_{0,1}, HG_{1,1}, HG_{2,1}, HG_{0,2}, HG_{1,2}, HG_{2,2}$

Multimode Vortex Beam



Conclusion

Applying quantum tomography to Shack-Hartmann data, the coherence matrix carrying complete information about the 3D intensity distribution of the signal can be reconstructed. As the method is a single-shot measurement and the Shack-Hartmann sensor is easy to build cost effective device, there is a great potential for implementing this technique in general 3-D imaging problems.

References

[1] B. Stoklasa, L. Motka, J. Rehacek, Z. Hradil and L.L. Sanchez-Soto, "Wavefront sensing reveals optical coherence," Nat. Commun. 5, 10.1038 (2014)

[2] Z. Hradil, J. Rehacek, and L.L. Sanchez-Soto, "Quantum Reconstruction of the Mutual Coherence Function," Phys. Rev. Lett. **105**, 010401 (2010).

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